

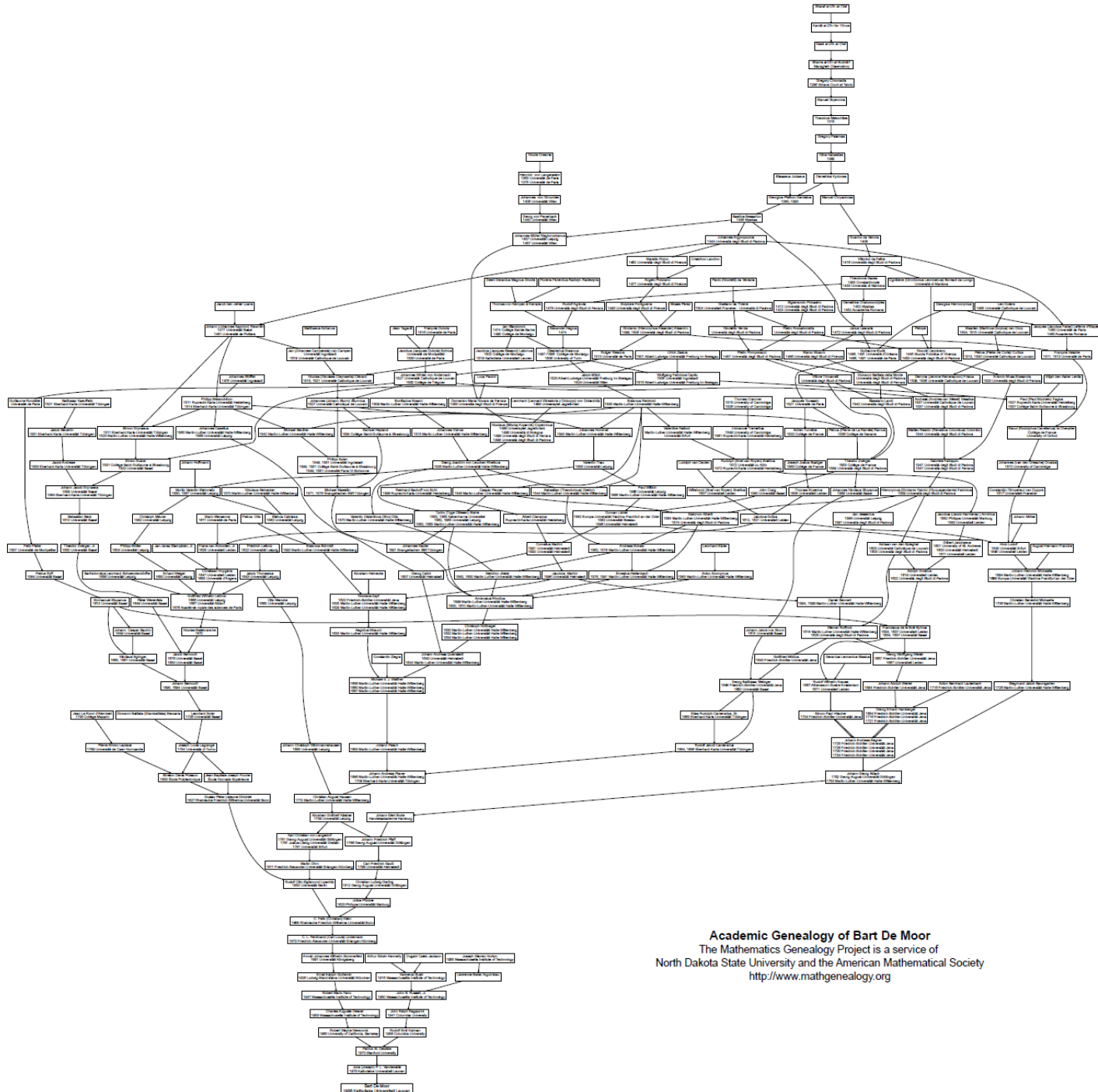
# Back to the Roots: Bart's legacy in the Quantum World

Frank Verstraete  
Ghent University  
Leuven 08/07/22

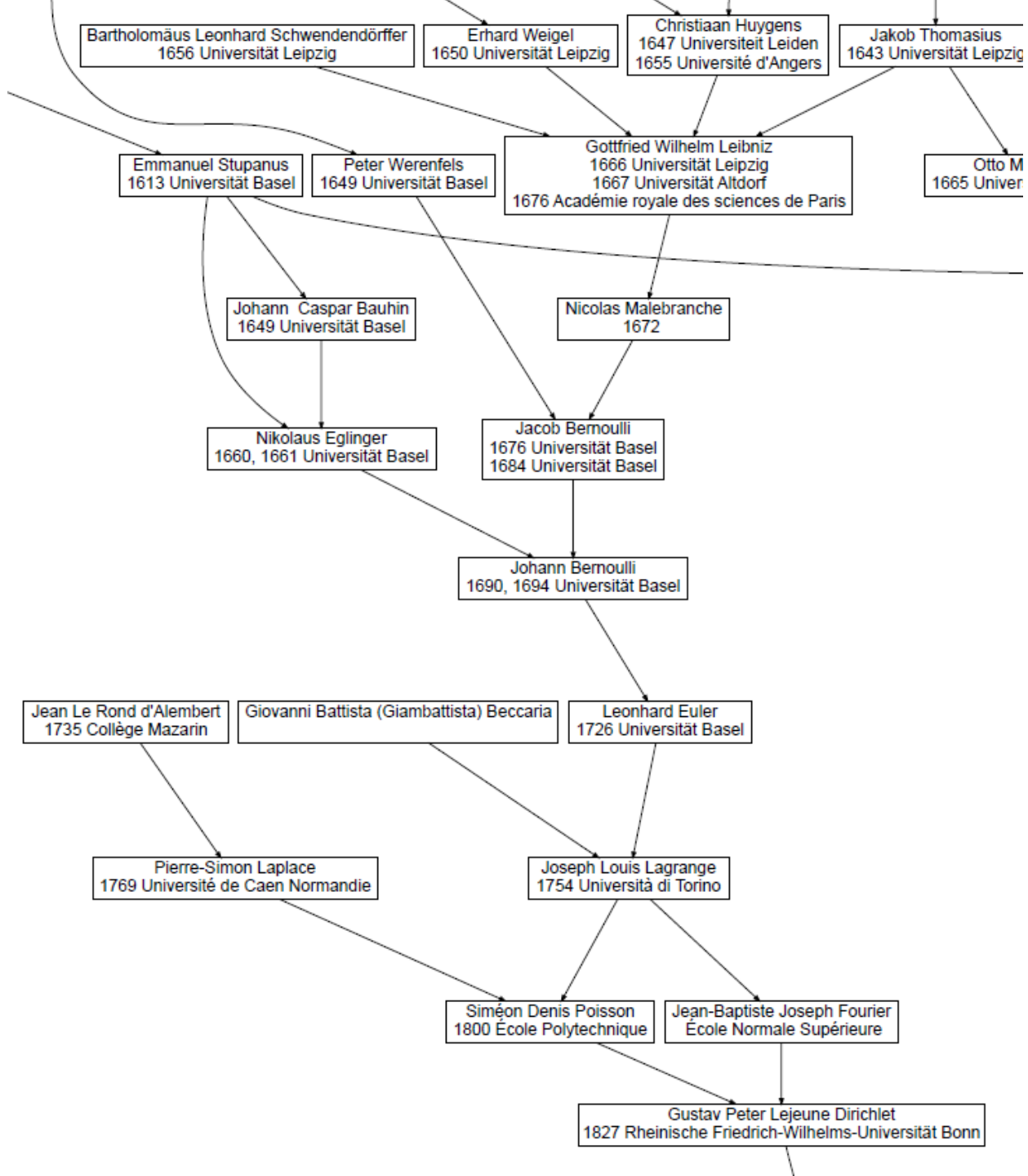
# It's the linear algebra, stupid!

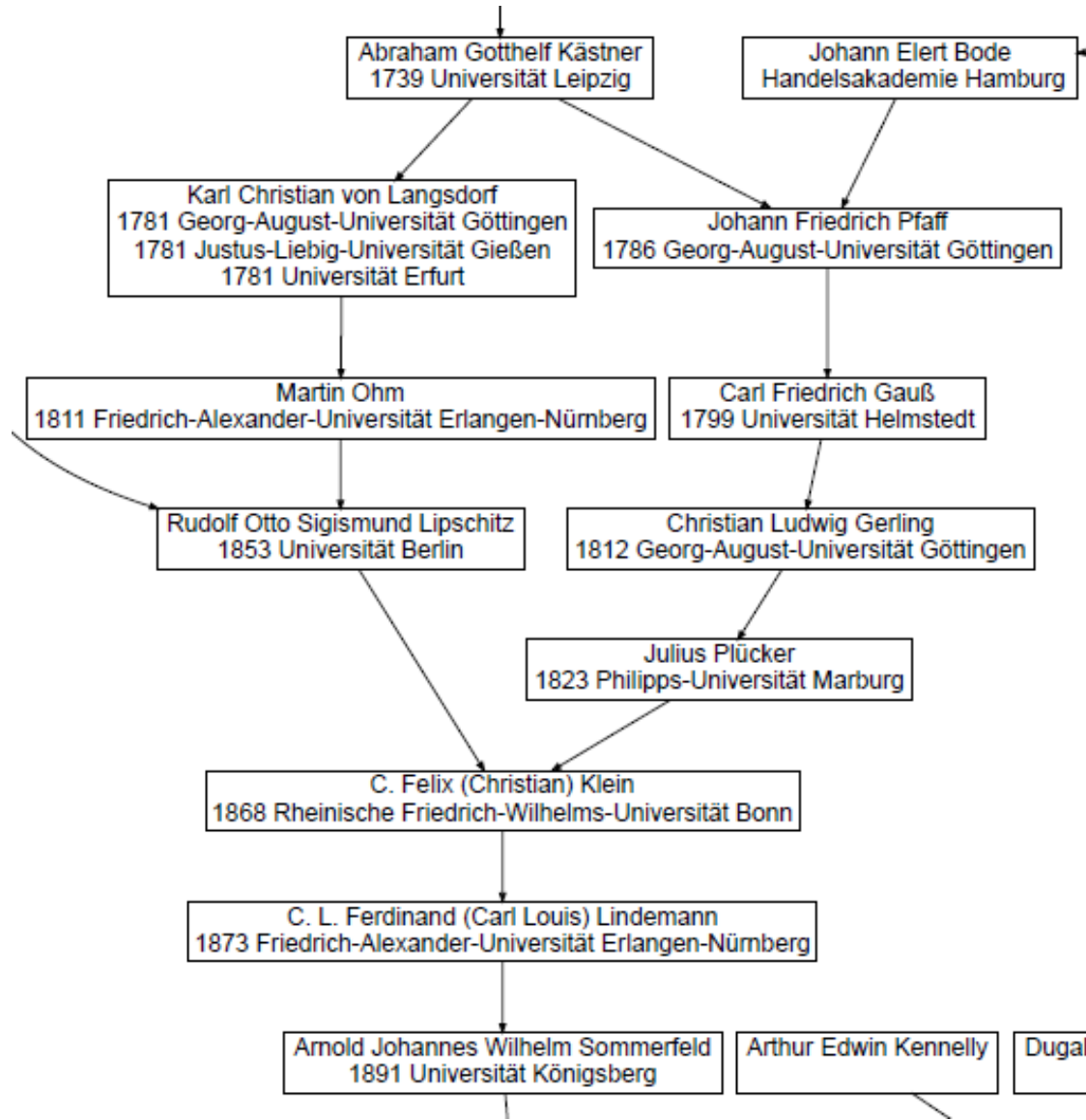
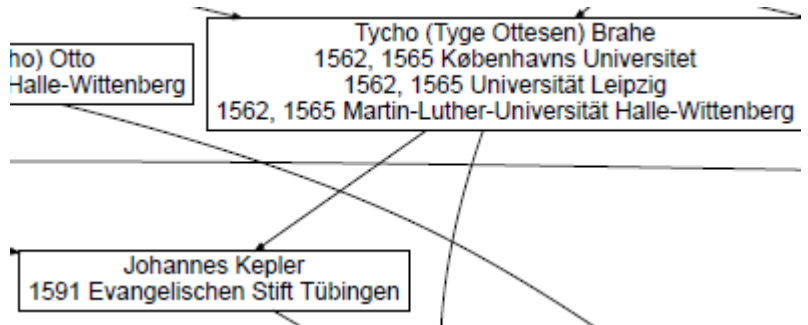
(Carville/Clinton/De Moor)

- I was raised in the group of Bart with the mantra “It’s the linear algebra, stupid”
- Took this advice to heart, and applied it to the then (1999) new field of quantum entanglement / quantum information / quantum computation
- I was - among many many others - extremely fortunate to have an advisor like Bart who believed in me – and gave me complete freedom in exploring my own interests
- The passion, enthusiasm and interests of Bart are unsurpassed – he is the ultimate post-Renaissance-man and role model - and therefore he had to understand quantum physics, the most successful theory of all matter



**Academic Genealogy of Bart De Moor**  
 The Mathematics Genealogy Project is a service of  
 North Dakota State University and the American Mathematical Society  
<http://www.mathgenealogy.org>





# Entanglement: what is it?

- Quantum states are described as vectors in a Hilbert space (superposition principle):

$$|\psi_A\rangle = \sum_i \psi_i |i\rangle$$

- Most relevant feature of Hilbert space for our purposes: the tensor product structure for many body systems

$$|\psi_{AB}\rangle = \sum_{ij} \psi_{ij} |i\rangle |j\rangle \qquad |\psi\rangle = \sum_{ijk\dots} X_{ijk\dots} |i\rangle |j\rangle \dots$$

- Any state which does not factorize with respect to the preferred basis (  $\psi_{ij} \neq \chi_i \cdot \xi_j$  ) is called entangled

### Four qubits can be entangled in nine different ways

[F.Verstraete](#), [J.Dehaene](#), [B.De Moor](#), [H.Vershelde](#) - Physical Review A, 2002 - APS


We consider a single copy of a pure four-partite state of qubits and investigate its behavior under the action of stochastic local quantum operations assisted by classical communication (...)

☆ Save  Cite Cited by 804 Related articles All 10 versions Web of Science: 472

### A comparison of the entanglement measures negativity and concurrence

[F.Verstraete](#), [K.Audenaert](#), [J.Dehaene](#)... - Journal of Physics A ..., 2001 - iopscience.iop.org

In this paper we investigate two different entanglement measures in the case of mixed states of two qubits. We prove that the negativity of a state can never exceed its concurrence and ...

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### Variational characterizations of separability and entanglement of formation

[K.Audenaert](#), [F.Verstraete](#), [B.De Moor](#) - Physical Review A, 2001 - APS

In this paper we develop a mathematical framework for the characterization of separability and entanglement of formation (EOF) of general bipartite states. These characterizations are ...

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### Normal forms and entanglement measures for multipartite quantum states

[F.Verstraete](#), [J.Dehaene](#), [B.De Moor](#) - Physical Review A, 2003 - APS

A general mathematical framework is presented to describe local equivalence classes of multipartite quantum states under the action of local unitary and local filtering operations. This ...

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### Maximally entangled mixed states of two qubits

[F.Verstraete](#), [K.Audenaert](#), [B.De Moor](#) - Physical Review A, 2001 - APS

We consider mixed states of two qubits and show under which global unitary operations their entanglement is maximized. This leads to a class of states that is a generalization of the ...

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### Lorentz singular-value decomposition and its applications to pure states of three qubits

[F.Verstraete](#), [J.Dehaene](#), [B.De Moor](#) - Physical Review A, 2002 - APS

All mixed states of two qubits can be brought into normal form by the action of local operations and classical communication operations of the kind  $\rho = (A \otimes B) \rho (A \otimes B)^\dagger$ . These normal ...

☆ Save  Cite Cited by 87 Related articles All 12 versions Web of Science: 49

### Local permutations of products of Bell states and entanglement distillation

[J.Dehaene](#), [M.Van den Nest](#), [B.De Moor](#), [F.Verstraete](#) - Physical Review A, 2003 - APS

We present different algorithms for mixed-state multicopy entanglement distillation for pairs of qubits. Our algorithms perform significantly better than the best-known algorithms. Better ...

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### On the geometry of entangled states

[F.Verstraete](#), [J.Dehaene](#), [B.De Moor](#) - journal of modern optics, 2002 - Taylor & Francis


The basic question that is addressed in this paper is finding the closest separable state for a given entangled state, measured with the Hilbert-Schmidt distance. While this problem is in ...

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### Local filtering operations on two qubits

[F.Verstraete](#), [J.Dehaene](#), [B.De Moor](#) - Physical Review A, 2001 - APS

We consider one single copy of a mixed state of two qubits and investigate how its entanglement changes under local quantum operations and classical communications (LQCC) of the ...

☆ Save  Cite Cited by 220 Related articles All 7 versions Web of Science: 139

# Four qubits can be entangled in nine different ways

F. Verstraete, J. Dehaene, B. De Moor, and H. Verschelde  
Phys. Rev. A **65**, 052112 – Published 25 April 2002

Article

References

Citing Articles (442)

PDF

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## ABSTRACT

We consider a single copy of a pure four-partite state of qubits and investigate its behavior under the action of stochastic local quantum operations assisted by classical communication (SLOCC). This leads to a complete classification of all different classes of pure states of four qubits. It is shown that there exist nine families of states corresponding to nine different ways of entangling four qubits. The states in the generic family give rise to Greenberger-Horne-Zeilinger-like entanglement. The other ones contain essentially two- or three-qubit entanglement distributed among the four parties. The concept of concurrence and 3-tangle is generalized to the case of mixed states of four qubits, giving rise to a seven-parameter family of entanglement monotones. Finally, the SLOCC operations maximizing all these entanglement monotones are derived, yielding the optimal single-copy distillation protocol.

Received 29 November 2001



# Workshop: Recognizing wildness in all its guises

- Workshop in Santa Fe, 2015

The workshop will focus on wildness in physics, math, and computer science. A classification problem is wild if it contains the problem of classifying the representations of any finite-dimensional algebra. **In physics, the perhaps more interesting example of wildness is the classification of entangled states up to local unitaries or local operations with classical communication (as much as possible - a complete classification in general is incredibly unlikely).** The geometry of the local unitary orbits are current frontier examples to try to understand in the **Geometric Complexity Theory program towards P vs NP**. In CS, this same problem is furthermore closely related to the complexity of Graph Isomorphism, which is neither known to be in P nor known to be NP-complete (a rare gem). These questions are also essentially the same as those arising at the intersection of representation theory and algebraic geometry that mathematicians have been interested in for decades.

Despite these connections, these three disciplines - physics, CS, and math - have largely been studying these problems from different angles and with different tools. The workshop will bring them together, with special attention paid to bridging the language barriers between these fields, in the hope that combining ideas and methods from different fields can lead to a better understanding in all of them.

$$|\psi'\rangle = A_1 \otimes A_2 \otimes A_3 \otimes A_4 |\psi\rangle$$

*Theorem 1.* Given a complex  $n \times n$  matrix  $R$ , then there always exist complex square orthogonal matrices  $O_1$  and  $O_2$  such that  $R' = O_1 R O_2$  is a unique direct sum of blocks of the form

(1)  $m \times m$  blocks of the form  $(\lambda_j I_m + S_m)$  being symmetric Jordan blocks (see, for example, [14] 4.4.9), and  $\lambda_j$  is a complex parameter (note that the case  $m = 1$  corresponds to the scalar case);

(2)  $m \times m$  blocks consisting of an upper left  $(m_1 + 1) \times m_1$  part being the matrix obtained by taking the odd rows and even columns of an  $(2m_1 + 1) \times (2m_1 + 1)$  symmetric Jordan block, and a lower right  $(m - m_1 - 1) \times (m - m_1)$  part being the transpose of the matrix obtained by taking the odd rows and even columns of a  $[2(m - m_1) - 1] \times [2(m - m_1) - 1]$  symmetric Jordan block.

$$G_{abcd} = \frac{a+d}{2} (|0000\rangle + |1111\rangle) + \frac{a-d}{2} (|0011\rangle + |1100\rangle) \\ + \frac{b+c}{2} (|0101\rangle + |1010\rangle) + \frac{b-c}{2} (|0110\rangle + |1001\rangle),$$

$$L_{abc_2} = \frac{a+b}{2} (|0000\rangle + |1111\rangle) + \frac{a-b}{2} (|0011\rangle + |1100\rangle) \\ + c (|0101\rangle + |1010\rangle) + |0110\rangle,$$

$$L_{a_2 b_2} = a (|0000\rangle + |1111\rangle) + b (|0101\rangle + |1010\rangle) + |0110\rangle \\ + |0011\rangle,$$

$$L_{ab_3} = a (|0000\rangle + |1111\rangle) + \frac{a+b}{2} (|0101\rangle + |1010\rangle) \\ + \frac{a-b}{2} (|0110\rangle + |1001\rangle) + \frac{i}{\sqrt{2}} (|0001\rangle + |0010\rangle) \\ + |0111\rangle + |1011\rangle),$$

$$L_{a_4} = a (|0000\rangle + |0101\rangle + |1010\rangle + |1111\rangle) + (i|0001\rangle \\ + |0110\rangle - i|1011\rangle),$$

$$L_{a_2 0_{3\oplus 1}} = a (|0000\rangle + |1111\rangle) + (|0011\rangle + |0101\rangle + |0110\rangle),$$

$$L_{0_{5\oplus 3}} = |0000\rangle + |0101\rangle + |1000\rangle + |1110\rangle,$$

$$L_{0_{7\oplus 1}} = |0000\rangle + |1011\rangle + |1101\rangle + |1110\rangle,$$

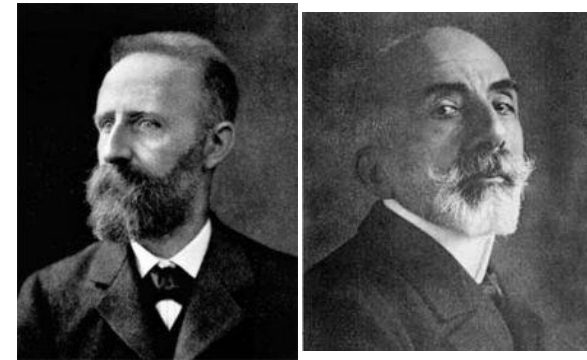
$$L_{0_{3\oplus 1} \overline{0}_{3\oplus 1}} = |0000\rangle + |0111\rangle.$$

# A complete set of covariants of the four qubit system

Emmanuel Briand, Jean-Gabriel Luque and Jean-Yves Thibon

Institut Gaspard Monge, Université de Marne-la-Vallée  
77454 Marne-la-Vallée cedex, France

We shall give here the first complete solution to the problem of describing the polynomial covariants, an investigation started by Le Paige as early as 1881 [7]. The theory was further advanced by C Segre in 1922 [16], using only geometric methods which led him close to a complete classification of the orbits. Such a complete classification was obtained only recently by Verstratete et al. [21], by exploiting the local isomorphism between  $SO_4$  and  $SL_2 \times SL_2$ , which permits a reduction of the problem to the classification of complex symmetric matrices up to orthogonal transformations.



## Constantin Le Paige

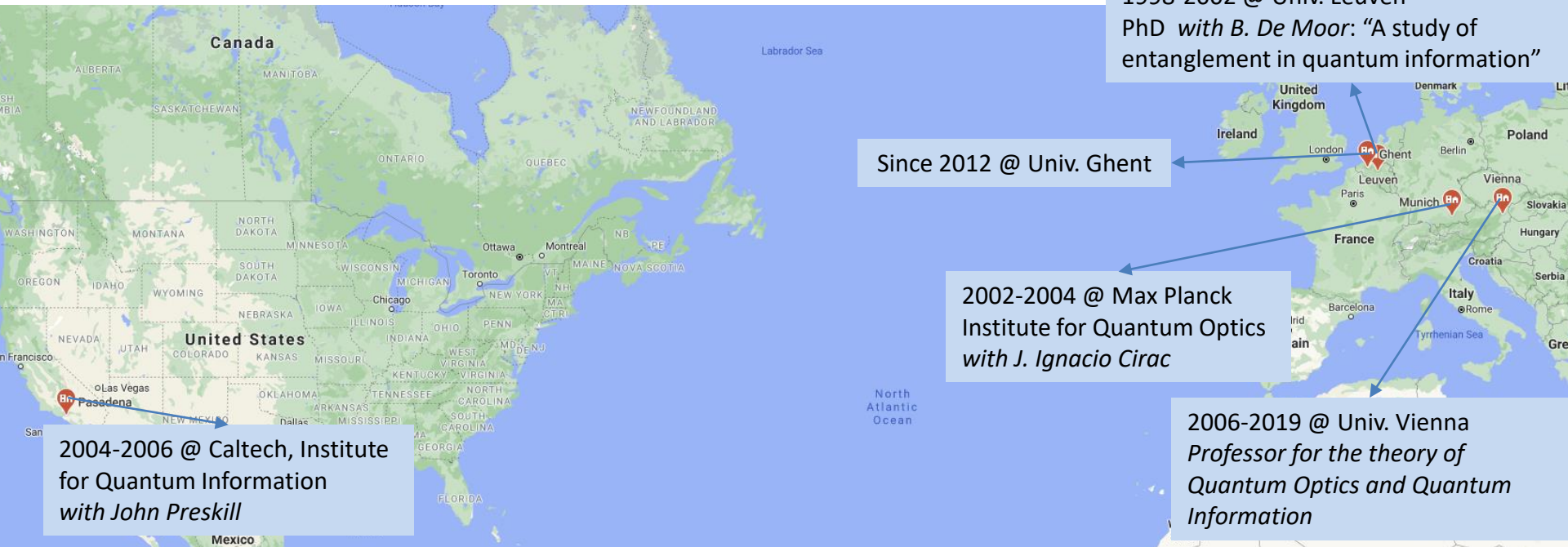
From Wikipedia, the free encyclopedia

**Constantin Marie Le Paige** (9 March 1852 – 26 January 1929) was a [Belgian mathematician](#).

Born in [Liège](#), [Belgium](#), Le Paige began studying [mathematics](#) in 1869 at the [University of Liège](#). After studying analysis under Professor [Eugène Charles Catalan](#), Le Paige became a professor at the Université de Liège in 1882.

While interested in [astronomy](#) and the [history of mathematics](#), Le Paige mainly worked on the theory of algebraic form, especially algebraic curves and surfaces and more particularly for his work on the construction of [cubic surfaces](#). Le Paige remained at the university until his retirement in 1922.

# My Trajectory

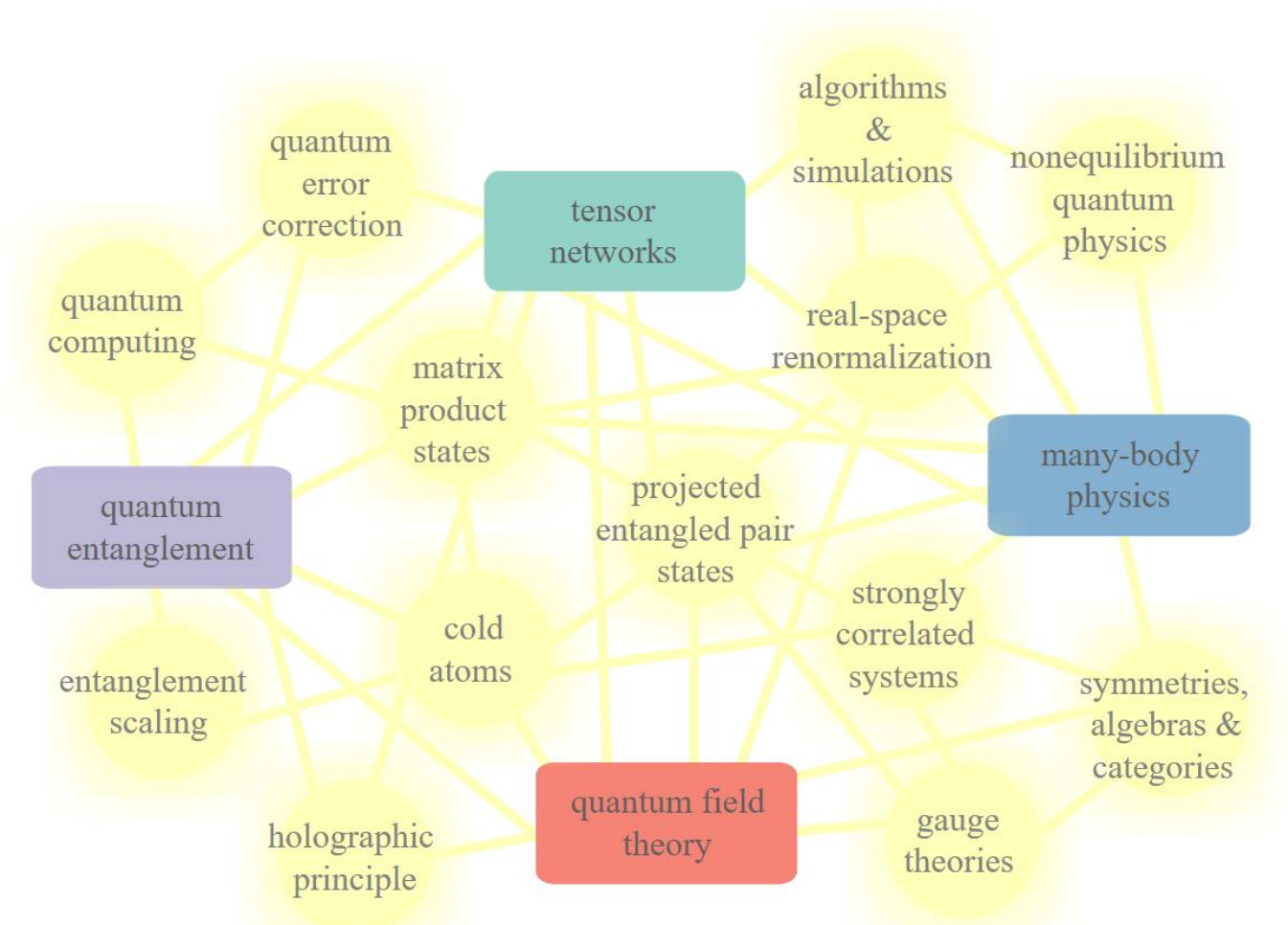


- Quantum Entanglement
  - Quantum Many-Body Physics
  - Quantum Computing
  - Strongly Correlated Systems
- } Tensor Networks

- Key contribution: realization that entanglement gives a new perspective on fundamental problems plaguing traditional approaches to the quantum many-body problem, and at the same time provides a completely new language for tackling it
  - Vocabulary consists of entangled qubits, syntax of tensor networks
  - Grand goal: reveal semantics of entanglement patterns

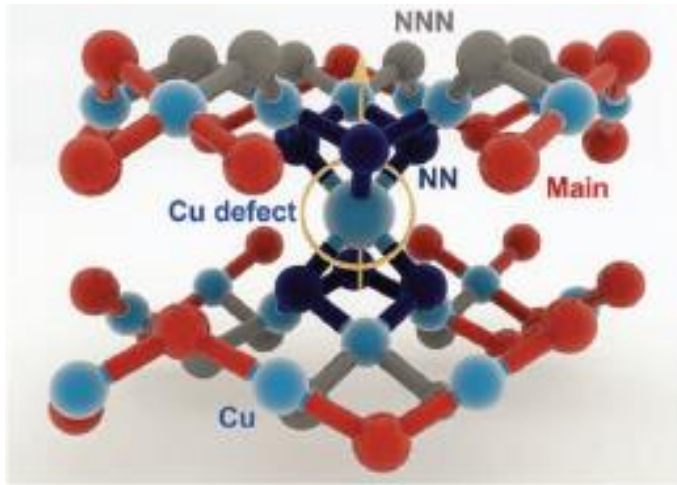
# Research

The QuantumGroup@UGent pursues a variety of research goals in the realm of theoretical physics; click on the items to find out more:



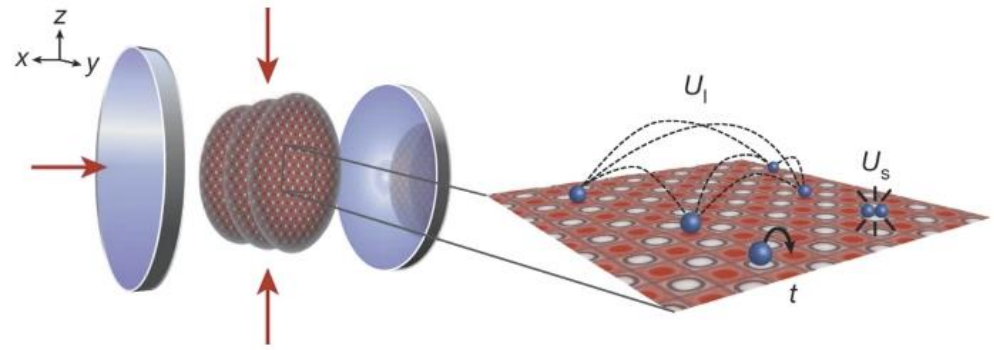
# The Quantum Many-Body Problem

- Problem: describe a large collection of interacting quantum particles
  - Is the central problem in quantum mechanics / physics since 1930's
    - Quantum chemistry, condensed matter, quantum field theory, ... : the whole is more than the sum of its parts
  - We know the (Schrödinger) equations, but cannot solve them: complexity increases exponentially in # particles
    - This scaling is both a curse (physics) and a blessing (computing)
  - Biggest breakthroughs in 20<sup>th</sup> century theoretical physics almost always consisted of the discovery of approximate methods for solving those equations
    - Hartree-Fock, Monte Carlo, Feynman's diagrammatic methods, Density Functional Theory, Renormalization Group, ...
    - All of those methods rely on the existence of a good fiducial noninteracting state, and those do not exist for e.g. the Hubbard model (high T<sub>c</sub> superconductivity)
  - Strongly correlated systems resist solutions using those traditional methods  
=> New methods needed!

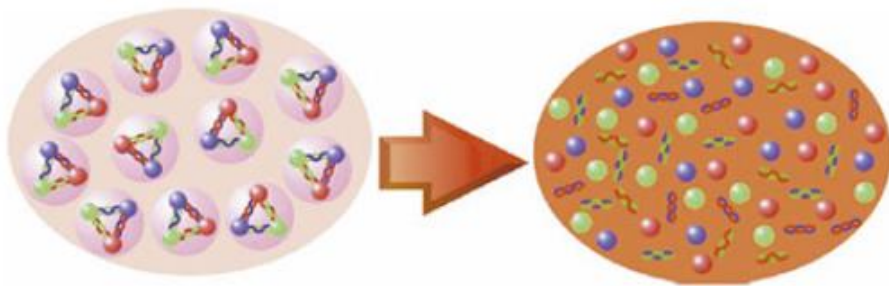


$\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$ .

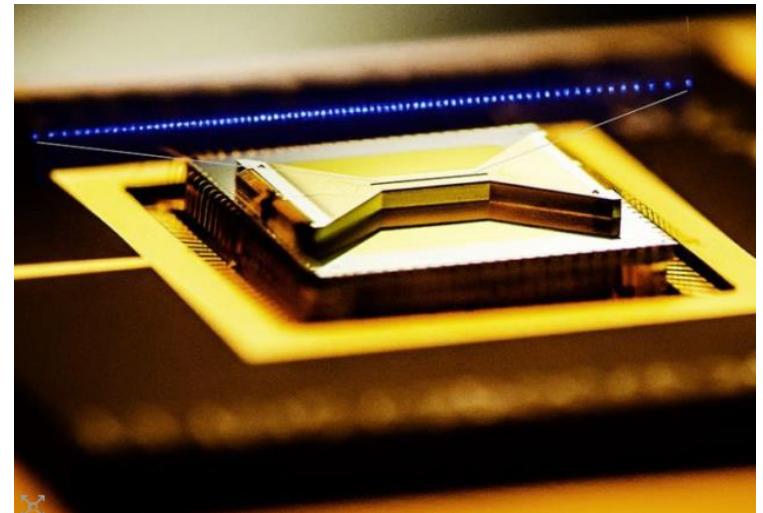
“Evidence for a gapped spin-liquid ground state in a kagome Heisenberg antiferromagnet”  
M. Fu et al., Science 350, 355 (2015)



“Quantum phases from competing short- and long-range interactions in an optical lattice”  
Landig et al, Nature 532, 476 (2016)



Quark-Gluon Plasma (taken from Riken)



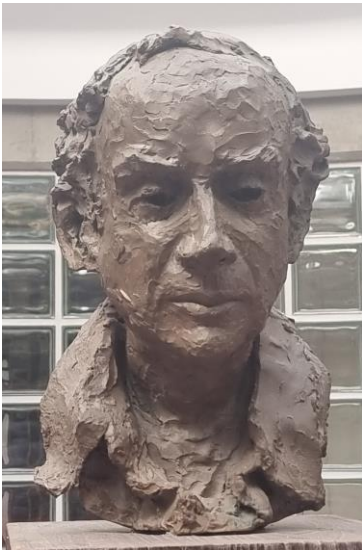
Ion trap quantum computer (taken from Rainer Blatt)

- Quantum physics of the 1920's:

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = \hat{H} |\Psi\rangle$$

- Hilbert space is endowed with a tensor product structure

$$|\psi\rangle = \sum_{i_1 i_2 \dots i_n} \psi_{i_1 i_2 i_3 \dots} |i_1\rangle |i_2\rangle$$

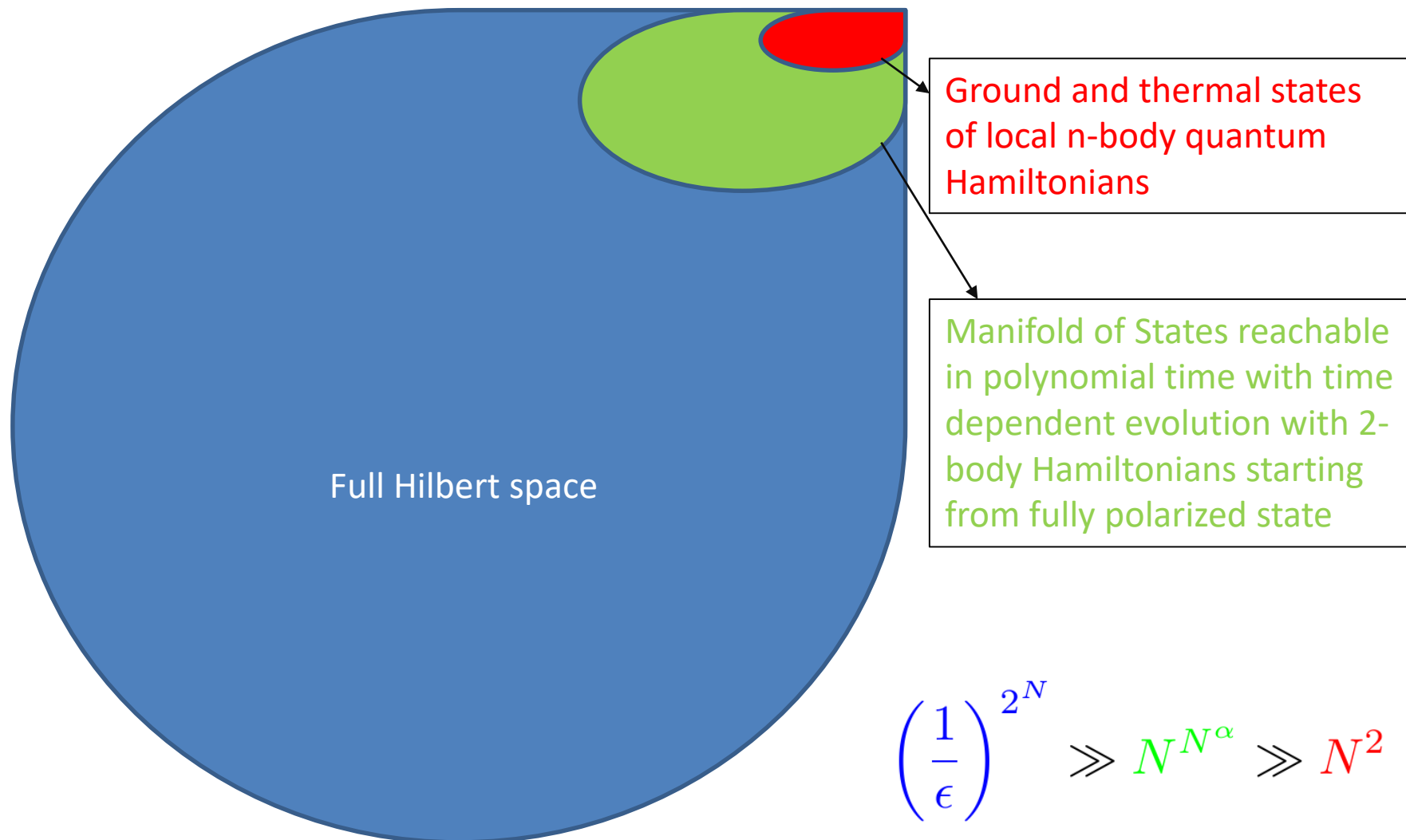


*“The underlying physical laws necessary for the mathematical theory of a large part of physics and the whole of chemistry are thus completely known, and **the difficulty is only that the exact application of these laws leads to equations much too complicated to be soluble.** It therefore becomes desirable that approximate practical methods of applying quantum mechanics should be developed, which can lead to an explanation of the main features of complex atomic systems without too much computation.” [Dirac ‘29]*



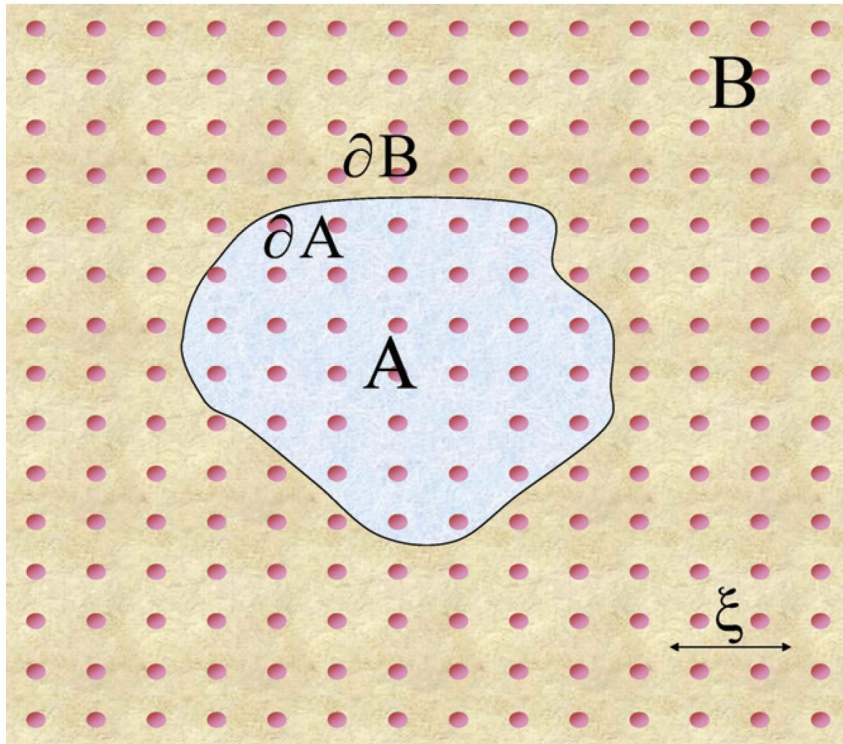
# Quantum Simulation of Time-Dependent Hamiltonians and the Convenient Illusion of Hilbert Space

David Poulin,<sup>1</sup> Angie Qarry,<sup>2,3</sup> Rolando Somma,<sup>4</sup> and Frank Verstraete<sup>2</sup>



# Area Laws for the entanglement entropy

- Ground and Gibbs states of interacting quantum many body Hamiltonians with **local interactions** have very peculiar properties
  - Area law for the entanglement entropy (ground states) or for mutual information (Gibbs states)
  - “explains” why physics is possible at all



Ground states:

$$S(\rho_A) = c \cdot \partial A$$

Srednicki '93; Hastings '07; ...

$$S(\rho_A) = \frac{c}{6} \cdot \log(A/\epsilon)$$

Holzhey, Larsen, Wilczek '94; ...

Gibbs states:

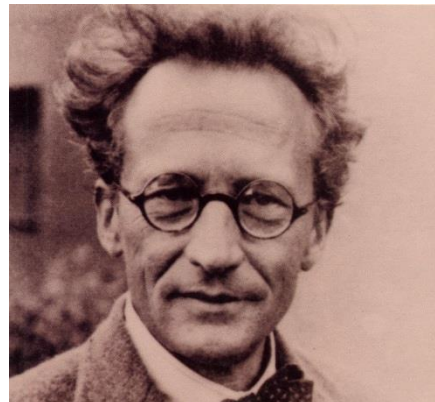
$$\begin{aligned} I(A, B) &= S(\rho_A) + S(\rho_B) - S(\rho_{AB}) \\ &= c \cdot \partial A \end{aligned}$$

Wolf, Hastings, Cirac, FV '08

# Quantum Physics in the 21<sup>st</sup> century



Entanglement



# Entanglement in an antiferromagnet



- Heisenberg antiferromagnet (spin  $\frac{1}{2}$ ):

$$\mathcal{H} = \sum_i \vec{S}_i \vec{S}_{i+1}$$

- What is the ground state?

– 2 spins:   $(|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)$

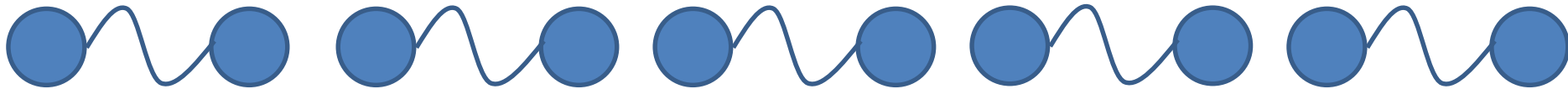
- 3 spins:

  $(|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)|\uparrow\rangle$        $-$        $|\uparrow\rangle(|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)$  

- Monogamy of entanglement: impossibility of sharing a singlet with two spin  $\frac{1}{2}$ 's. Emerging phenomena arise because of this competition.

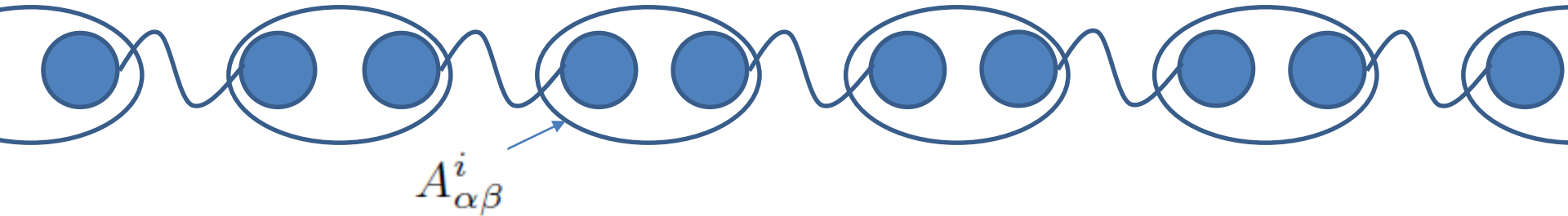
# Tensor Calculus for Quantum Spin Chains

- Systematic way of creating states which have extremal local marginals but keep translational invariance: *matrix product states (MPS)*



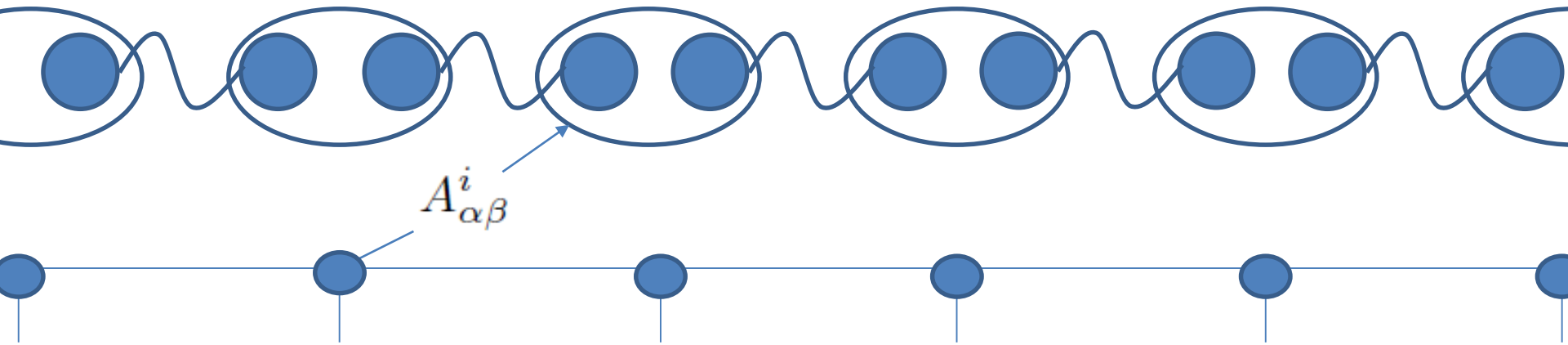
# Tensor Calculus for Quantum Spin Chains

- Systematic way of creating states which have extremal local marginals but keep translational invariance: **matrix product states (MPS)**



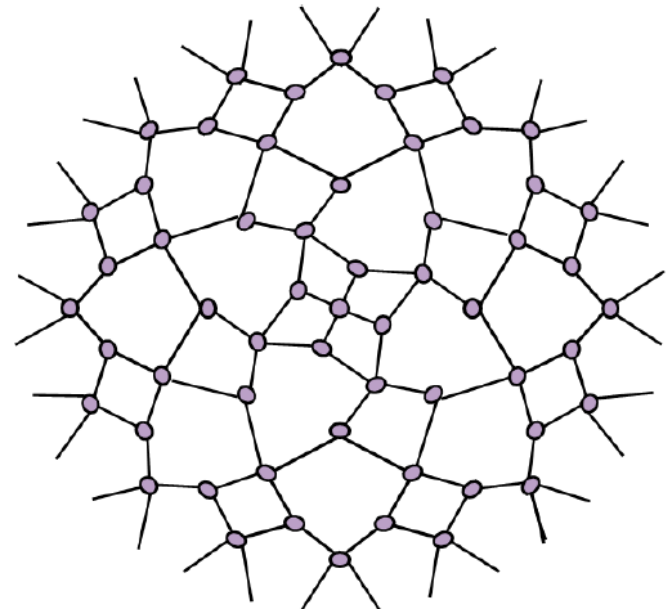
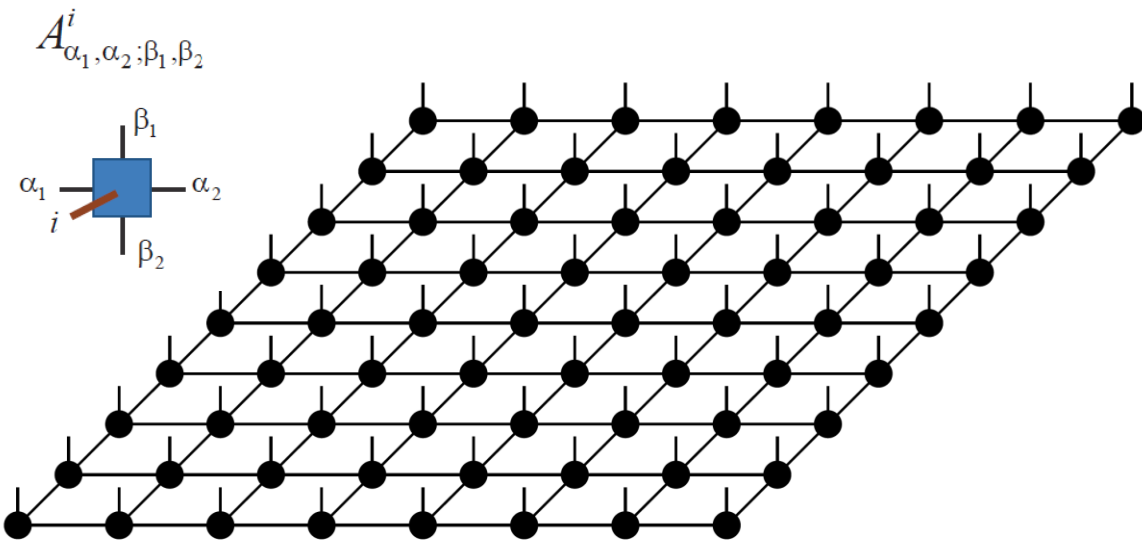
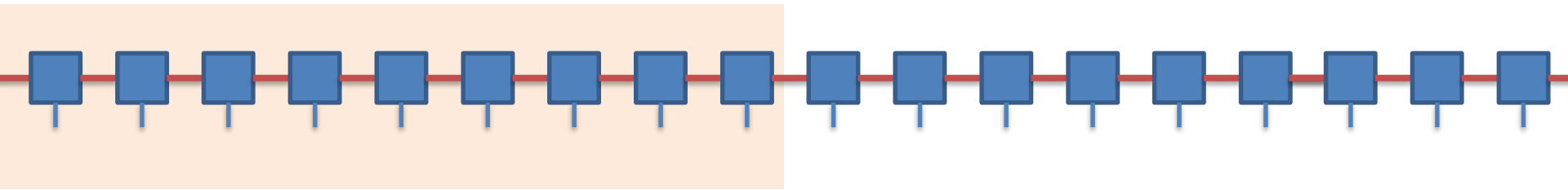
# Tensors for Quantum Spin Chains

- Systematic way of creating states which have extremal local marginals but keep translational invariance: **matrix product states (MPS)**



$$|\psi\rangle = \sum_{i_1 i_2 i_3 \dots} \text{Tr} (A^{i_1} A^{i_2} A^{i_3} \dots) |i_1\rangle |i_2\rangle |i_3\rangle \dots$$

# Entanglement as building block of matter: Quantum Tensor Networks



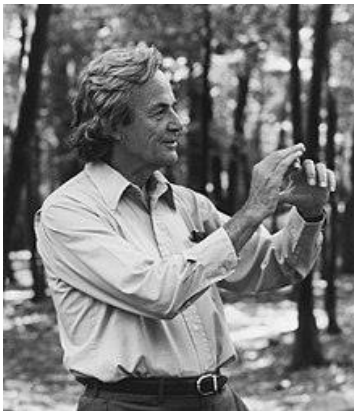


- ***tensor networks : crucial concepts***
  - Tensors model the entanglement structure: modelling correlations makes much more sense than modelling wavefunction directly
    - Tensors dictate the ***entanglement patterns***
  - Tensor networks can be efficiently contracted due to ***holographic property***: map quantum 3D  $\rightarrow$  2D  $\rightarrow$  1D  $\rightarrow$  0D problems, and this can be done efficiently due to area laws
  - States are defined in thermodynamic limit; *finite size scaling* is replaced by ***finite entanglement scaling***
  - *Local tensor contains all global information* about quantum many body state
    - different phases of matter can be distinguished by ***symmetries of those local tensors***, including topological phases
    - Tensor networks provide a natural way of dealing with gauge theories: enforcing symmetries

# Feynman's dream

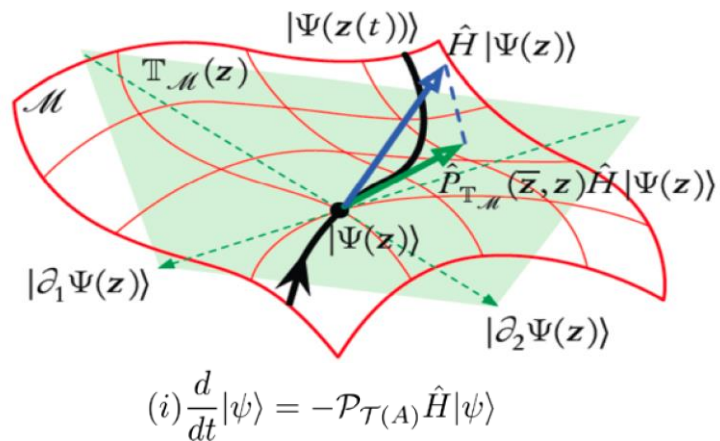
“Now, in field theory, what’s going on over here and what’s going on over there and all over space is more or less the same. **Why do we have to keep track in our functional of all things going on over there while we are looking at the things that are going on over here?** ... It’s really quite insane actually: we are trying to find the energy by taking the expectation of an operator which is located here and we present ourselves with a functional which is dependent on everything all over the map. That’s something wrong. **Maybe there is some way to surround the object, or the region where we want to calculate things, by a surface and describe what things are coming in across the surface. It tells us everything that’s going on outside.**

*I’m talking about a new kind of idea but that’s the kind of stuff we shouldn’t talk about at a talk, that’s the kind of stuff you should actually do!”*

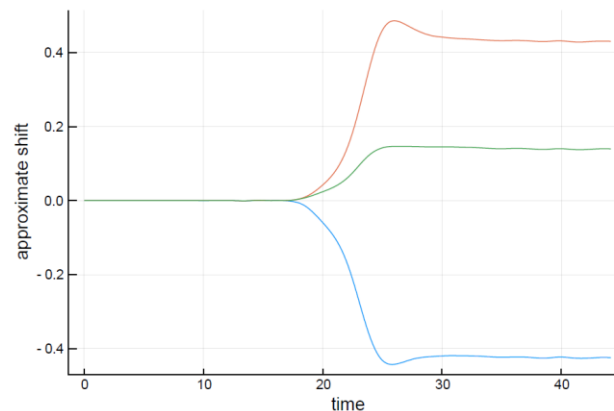


“Difficulties in Applying the Variational Principle to Quantum Field Theories”, Wangerooge 1987, Proceedings,  
*Variational calculations in quantum field theory*

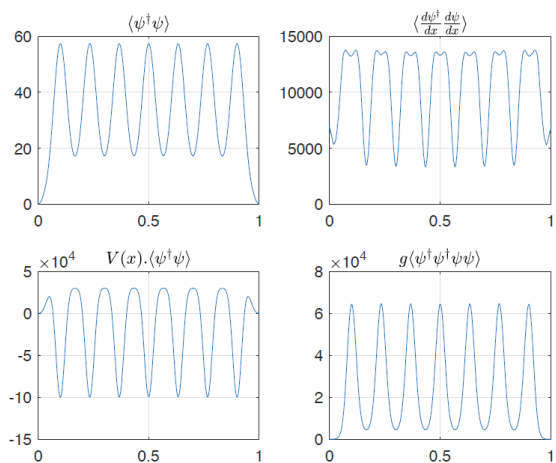
# MPS: computational aspects



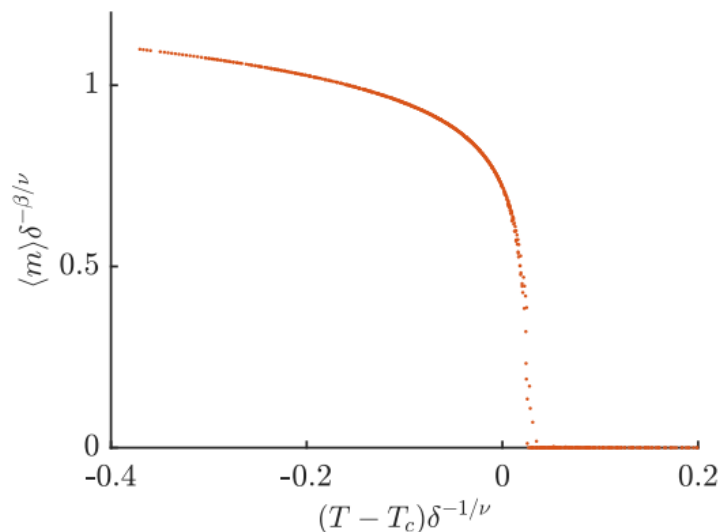
Time-Dependent Variational Principle for MPS  
PRL 107, 70601 (2011)



Scattering of magnons in spin 1 Heisenberg model  
PRR 3, 013078 (2021)

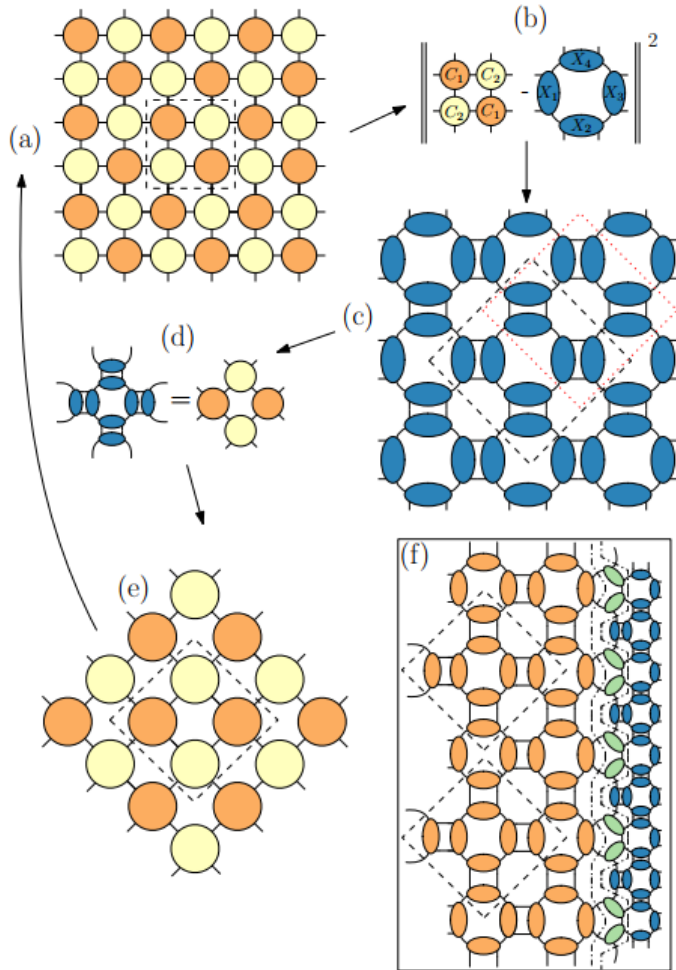


Interacting bosons with OBC:  
cMPS and quantum Gross-Pitaevskii  
PRL 128, 020501 (2022)

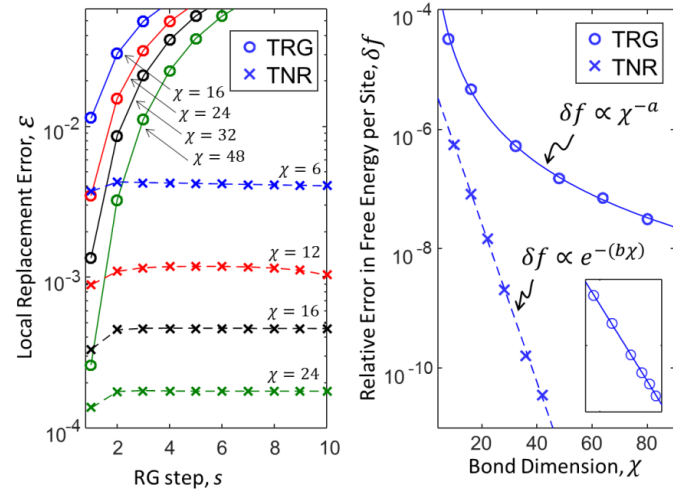


Finite Entanglement Scaling for critical Potts model  
PRL 123, 250604 (2019)

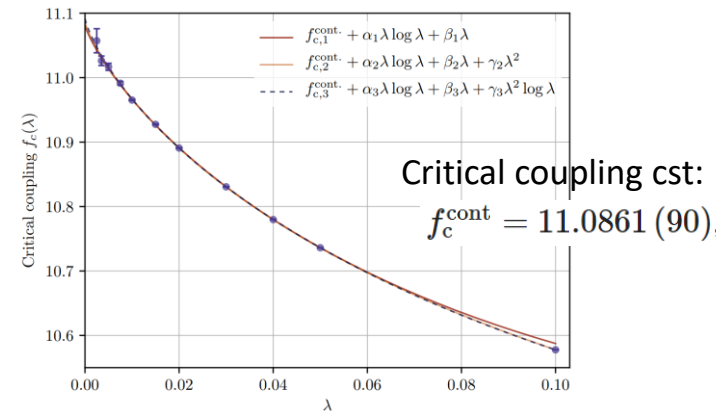
# Multiscale Entanglement Renormalization Ansatz (MERA), Tensor Renormalization Group (TRG), Tensor Network Renormalization (TNR)



Renormalization group flows of Hamiltonians using tensor networks  
PRL 118, 250602 (2017)



Tensor Network Renormalization Evenly, Vidal PRL 115, 180405 (2015)

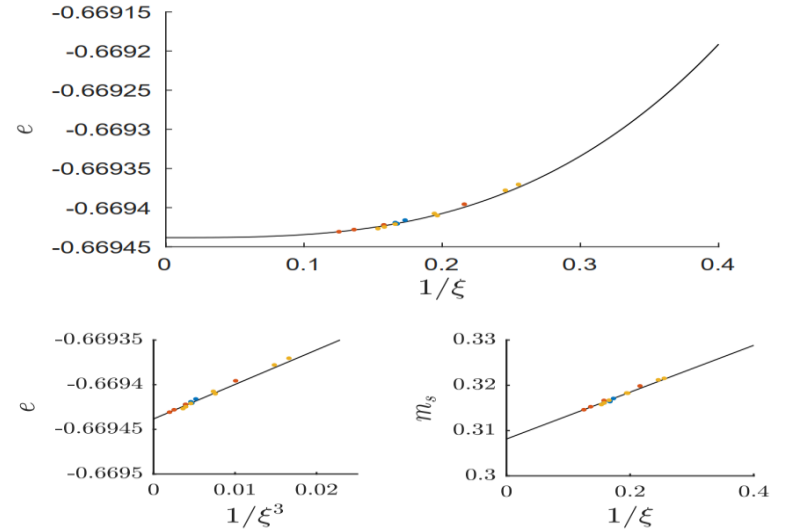


Computing the renormalization group flow of two-dimensional  $\Phi^4$  theory with tensor networks  
Delcamp, Tilloy PRR 2, 033278 (2020)

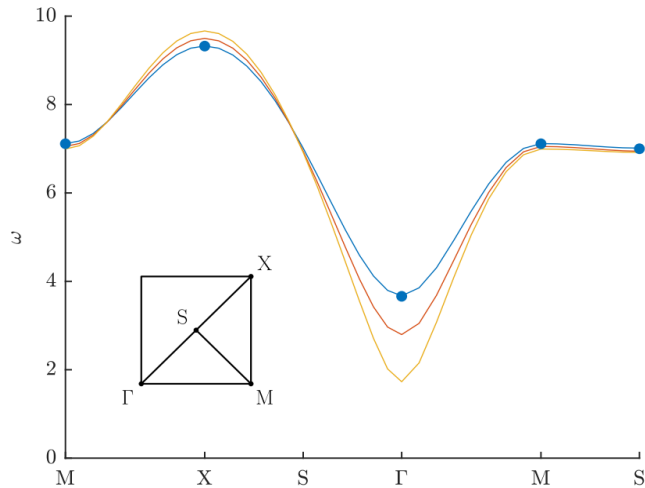
# Projected Entangled Pair States (PEPS)

Pauling <sup>13</sup>	mean field	1.5
Nagle <sup>15</sup>	series expansion	1.50685(15)
Berg et.al. <sup>19</sup>	multicanonical	1.507117(35)
Herrero et.al. <sup>16</sup>	num. integration	1.50786(12)
Kolafa <sup>17</sup>	num. integration	1.5074660(36)
PEPS	$D = 2$	1.50735
	$D = 3$	1.507451
	$D = 4$	1.507456

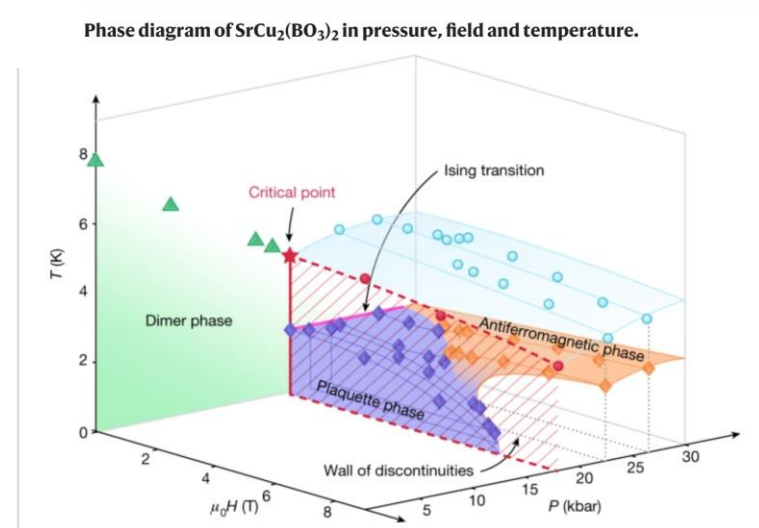
Spin Ice: residual entropy  
PRE 98, 042145 (2018)



Entanglement scaling for 2D Heisenberg model  
arXiv:2102.03143



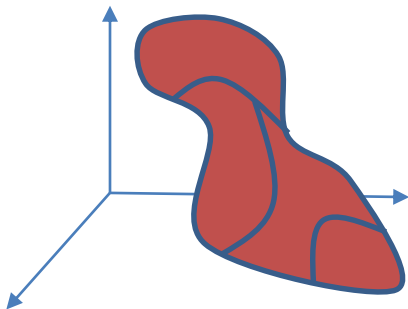
Dispersion relation for 2D Ising with transverse field  
PRB 99, 165121 (2019)



Physical realization of Shastry-Sutherland model  
Mila et al., Nature 592, 370 (2021)

# It's all about symmetries

- Anderson '72: “*physics is an applied form of group theory*”
- Essential paradigm: detect the **global** features of a system through its entanglement degrees of freedom / **local** tensor network description
  - The symmetries of the local tensor in the tensor network reveal the emerging properties of the system
  - Even for the case of topological order: *local* order parameters arise in the form of different representations of groups / fusion algebras
- This has led to the program of classifying phases of matter



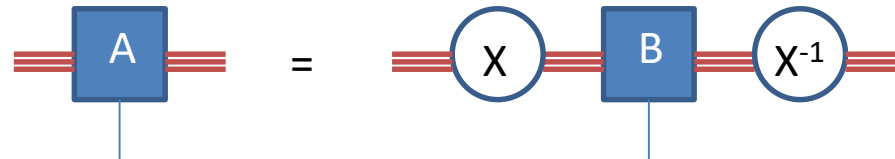
- QIT tools enable the correct definition of phases of matter: in terms of low-depth quantum circuits

Bravyi, Hastings, FV PRL 97, 050401 (2006)

# Tensor Networks: theoretical aspects

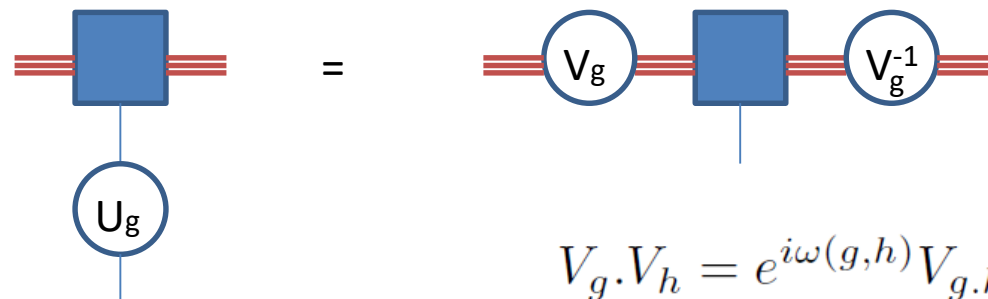
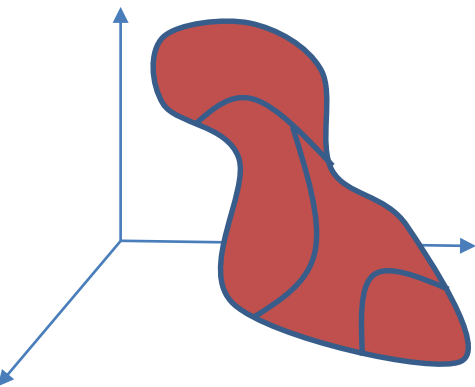
- MPS: fundamental theorem

$$|\psi(A^i)\rangle = |\psi(B^i)\rangle \Leftrightarrow \exists X : A^i = X B^i X^{-1}$$



Cirac, Perez-Garcia, Schuch, FV '16

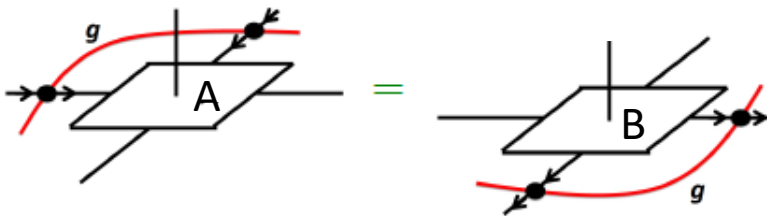
- Classification of 1D SPT phases



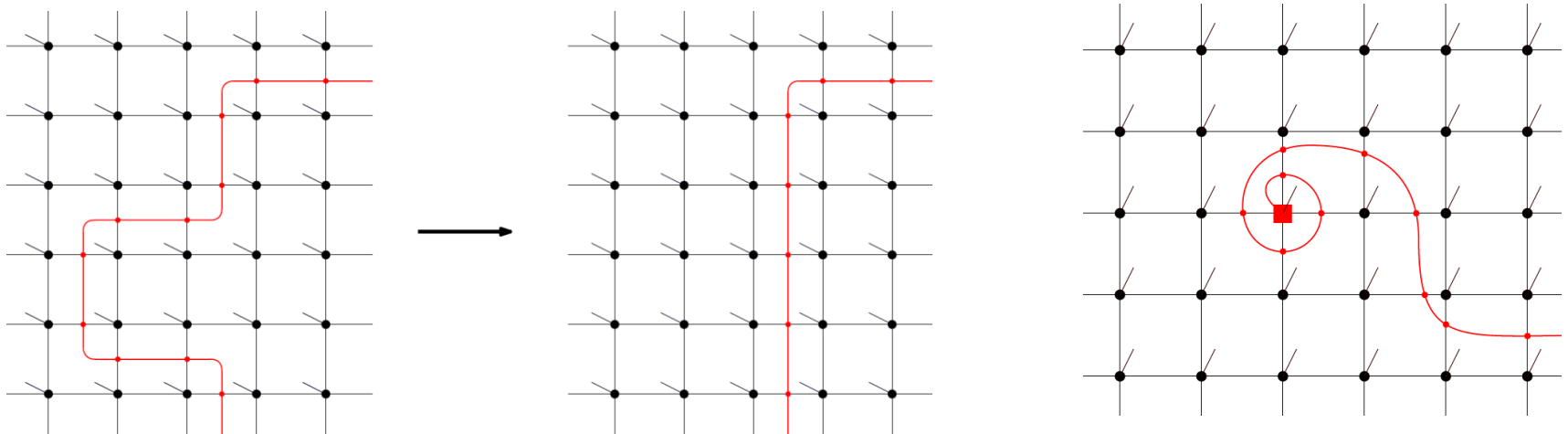
Perez-Garcia, Wolf, Sanz, FV, Cirac '08

Pollmann, Turner, Berg, Oshikawa '10; Chen, Gu, Wen '11; Schuch, Perez-Garcia, Cirac '11

- PEPS: fundamental theorem

$$|\psi(A^i)\rangle = |\psi(B^i)\rangle \iff$$


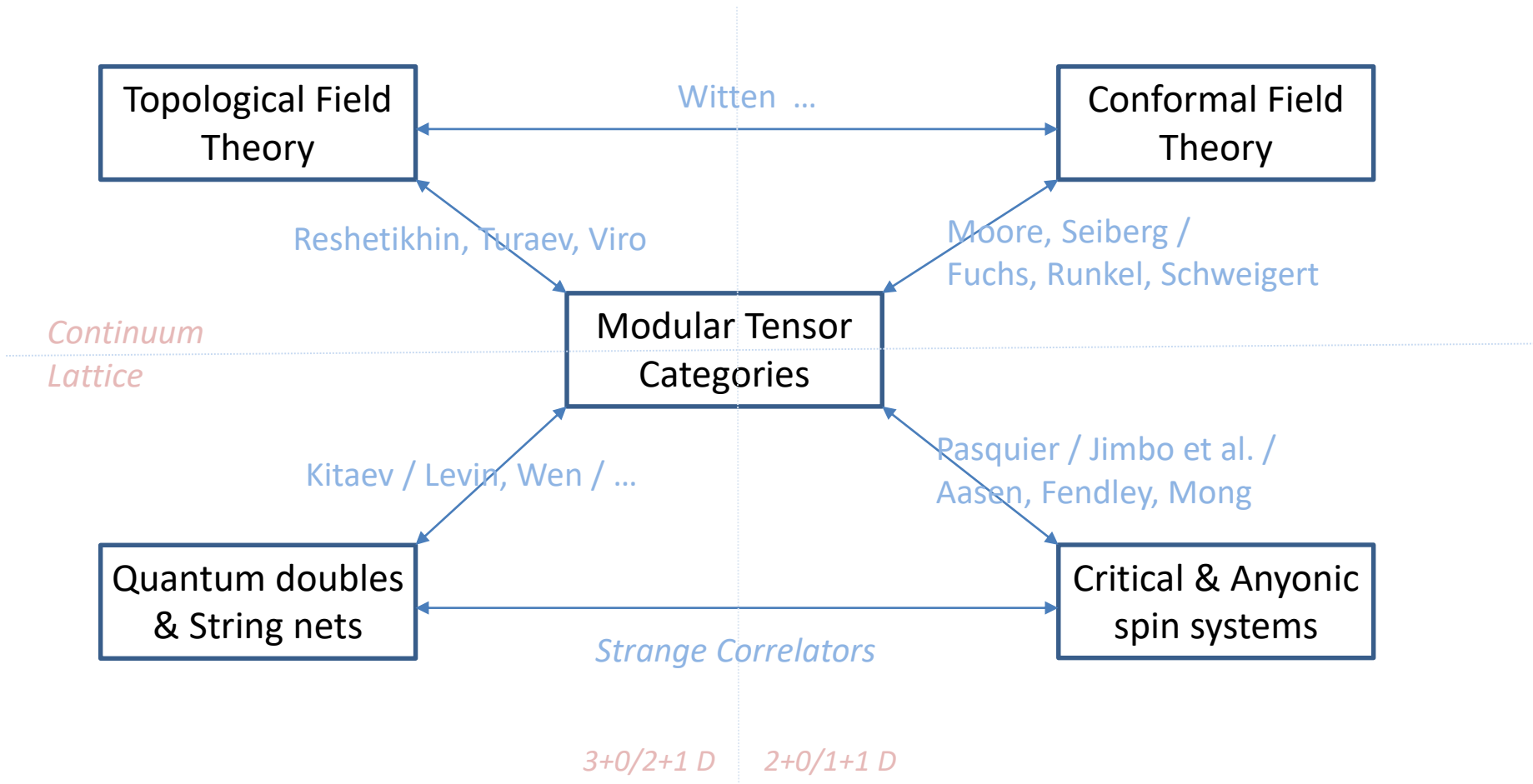
- Topological order: Levin-Wen string nets / Turaev-Viro state sums (TFT)



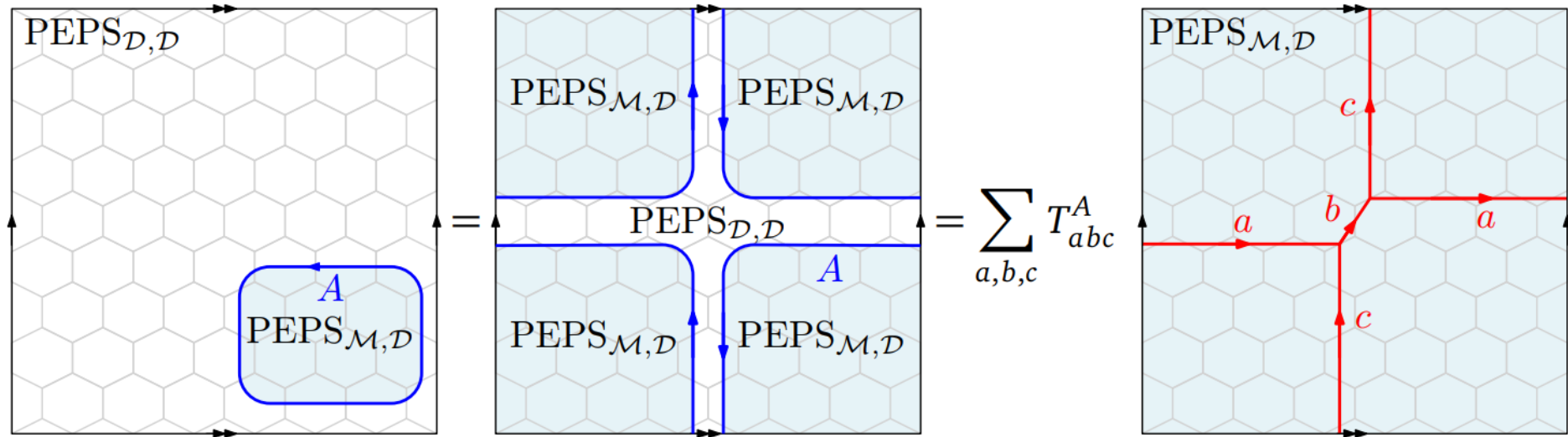
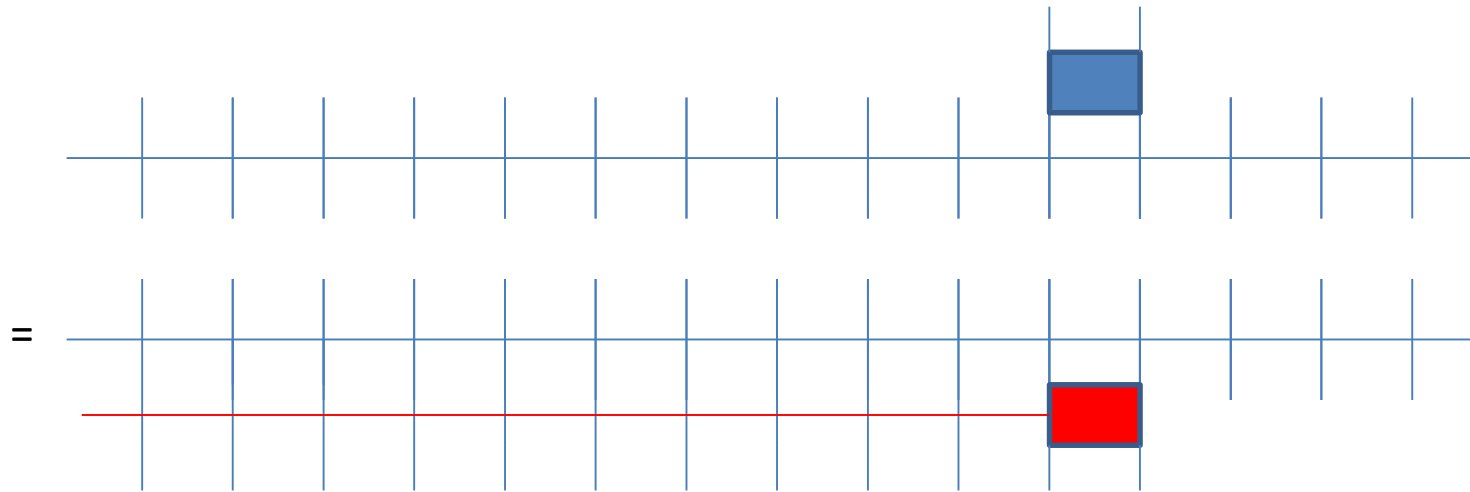
$$O_a \cdot O_b = \sum_c N_{ab}^c O_c$$



# Topological / Categorical symmetries: TFT & CFT



# Dualities and MPO intertwiners: bimodule categories

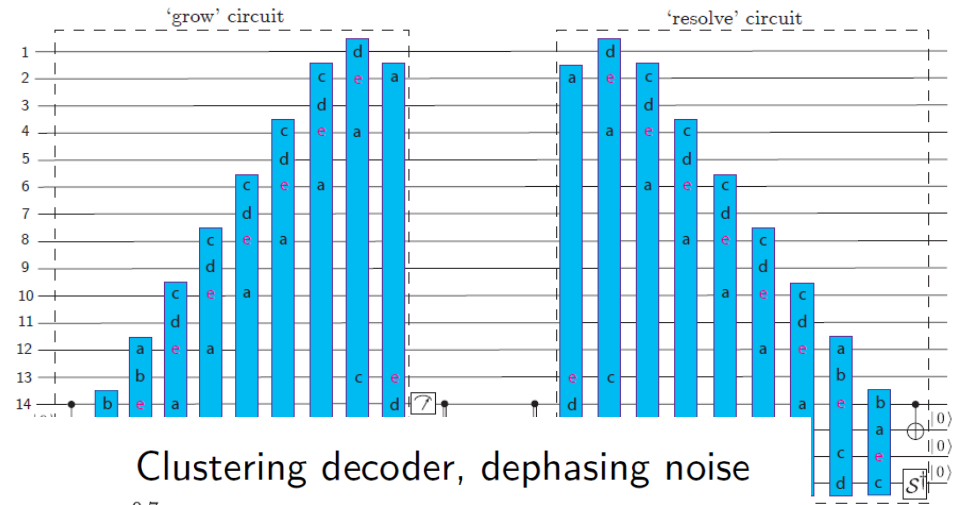
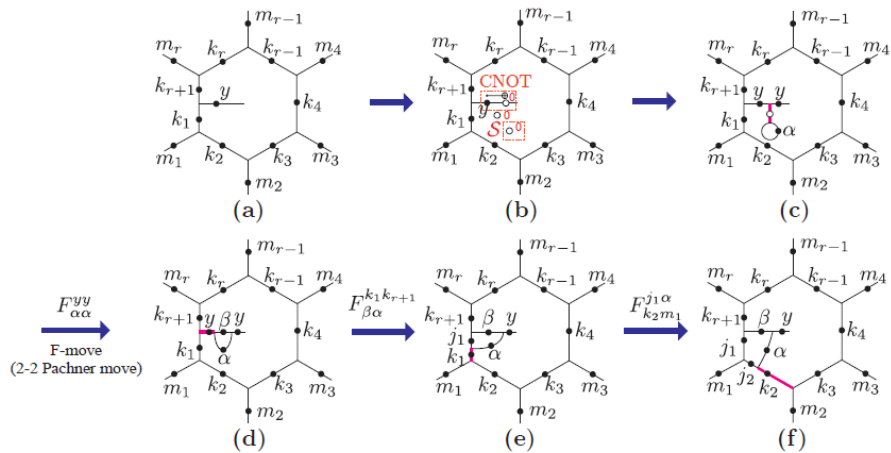


Category-theoretic recipe for dualities in one-dimensional quantum lattice models

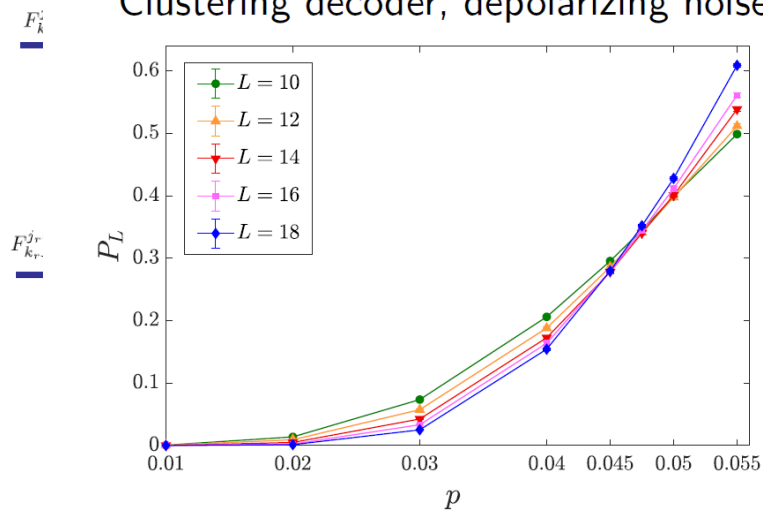
Lootens, Delcamp, Ortiz, FV arXiv:2112.09091



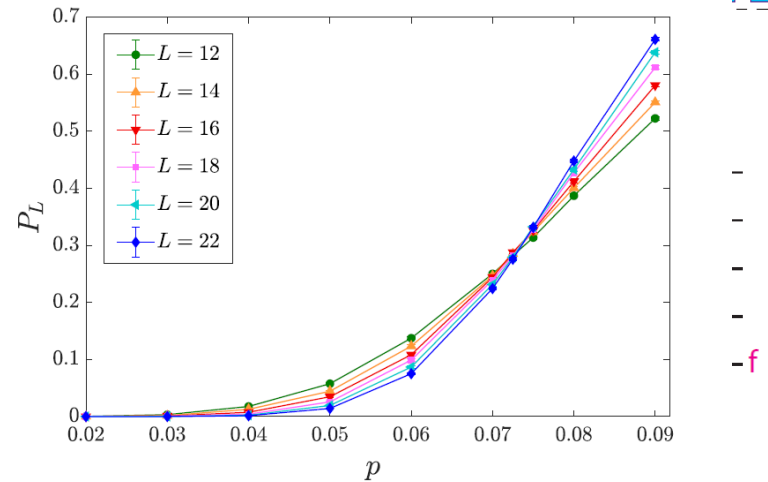
# Application: Fibonacci Turaev-Viro error correcting code



Clustering decoder, depolarizing noise



Clustering decoder, dephasing noise



# “the unreasonable effectiveness of ~~tensor networks~~ in physics”

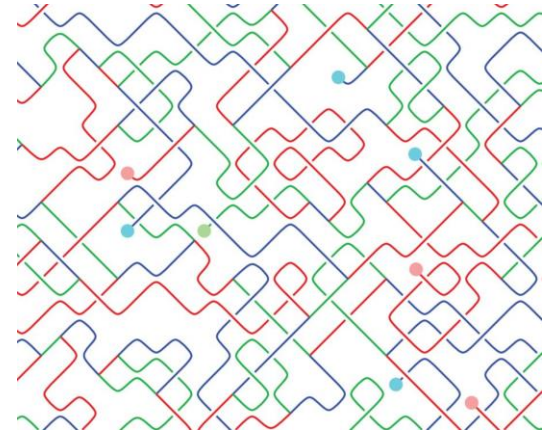
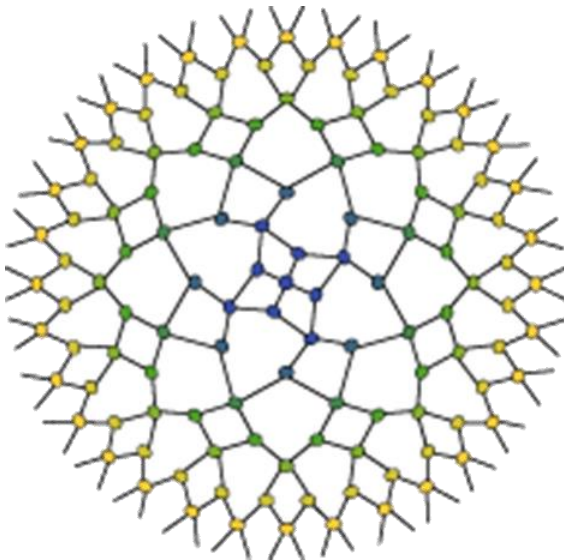
*Linear Algebra*

Real-Space Renormalization Group

Symmetries and dualities

Tensor Networks

Quantum Field Theory  
Statistical Mechanics



- High-Energy Physics
- Condensed Matter
- Cold Atoms

